

LABORATORY AND ASTRONOMICAL DETECTION OF THE NEGATIVE MOLECULAR ION C_3N^- P. THADDEUS,¹ C. A. GOTTLIEB,¹ H. GUPTA,^{1,2} S. BRÜNKEN,¹ M. C. MCCARTHY,¹
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ABSTRACT

The negative molecular ion C_3N^- has been detected at millimeter wavelengths in a low-pressure laboratory discharge, and then with frequencies derived from the laboratory data in the molecular envelope of IRC+10216. Spectroscopic constants derived from laboratory measurements of 12 transitions between 97 and 378 GHz allow the rotational spectrum to be calculated well into the submillimeter-wave band to 0.03 km s^{-1} or better in equivalent radial velocity. Four transitions of C_3N^- were detected in IRC+10216 with the IRAM 30 m telescope at precisely the frequencies calculated from the laboratory measurements. The column density of C_3N^- is 0.5% that of C_3N , or approximately 20 times greater than that of C_4H^- relative to C_4H . The C_3N^- abundance in IRC+10216 is compared with a chemical model calculation by Petrie & Herbst. An upper limit in TMC-1 for C_3N^- relative to C_3N ($<0.8\%$) and a limit for C_4H^- relative to C_4H ($<0.004\%$) that is 5 times lower than that found in IRC+10216, were obtained from observations with the NRAO 100 m Green Bank Telescope (GBT). The fairly high concentration of C_3N^- achieved in the laboratory implies that other molecular anions containing the CN group may be within reach.

Subject headings: ISM: molecules — line: identification — molecular data — molecular processes — radio lines: ISM

1. INTRODUCTION

The carbon chain radicals C_3N and C_4H , similar linear molecules with open shell $^2\Sigma$ electronic ground states and spin doubling and hyperfine structure in their rotational spectra, were identified 30 years ago prior to laboratory detection in the expanding molecular shell of the well-known evolved carbon star IRC+10216 (Guélin & Thaddeus 1977; Guélin et al. 1978). In the absence of laboratory data, the identifications were based on the rotational constants derived from the observations, which were within about 1% of those predicted, and on the well resolved spin doubling, which was about that expected. C_3N was identified first, and the subsequent observation of the predicted C_4H was evidence that the assignment was correct. Several years later laboratory identification conclusively confirmed both assignments (Gottlieb et al. 1983).

Positive molecular ions have long been known in the interstellar gas, starting with CH^+ , one of the first three interstellar molecules (Douglas & Herzberg 1941), and 14 are now known. Negative molecular ions, however, have only been known for just over one year; the first, C_6H^- , was shown on the basis of laboratory measurements (McCarthy et al. 2006) to be the carrier of an unidentified series of lines in IRC+10216 (Kawaguchi et al. 1995). C_4H^- was identified in a gas discharge in this laboratory in 2006 October (Gupta et al. 2007), and with rest frequencies in

hand was then quickly found in the millimeter-wave spectrum of IRC+10216 with the IRAM 30 m telescope (Cernicharo et al. 2007). Searches both in the laboratory and in space for the iso-electronic C_3N^- were then undertaken, and both have now succeeded. It is the purpose of the present paper to describe the laboratory and astronomical identification of this new negative ion.

2. LABORATORY AND ASTRONOMICAL OBSERVATIONS

2.1. Laboratory

We previously detected C_4H^- in both the millimeter-wave band with the free space spectrometer in this laboratory, and in the centimeter-wave band with the Fourier-transform microwave spectrometer, but so far we have only been able to observe C_3N^- with the free space instrument. The reason for this difference is unclear but, as discussed below, it may be fundamental, the result of the quite different discharge conditions employed in the two instruments. Twelve rotational lines between 97 and 378 GHz have been measured; Table 1 summarizes the data. The best-fit constants B and D in Table 2 reproduce the measured frequencies to an rms of 19 kHz, which is comparable to the measurement uncertainties; they allow the rotational spectrum to as high as 500 GHz to be calculated to 0.03 km s^{-1} in radial velocity, which is adequate for deep spectral line searches in nearly all astronomical sources. (The sixth-order centrifugal distortion constant H has not been determined, but to judge from the measured sixth-order term of HC_3N [Thorwirth et al. 2000], its effect on line frequencies even as high as 500 GHz is almost negligible.)

Our first attempts to find C_3N^- were extremely frustrating, and it was only after two surveys failed to yield evidence for this ion that a third independent search under different discharge conditions met with success. Each survey covered about 1.2 GHz of the millimeter-wave spectrum, and required about 10 hr of observation. The crucial parameter turned out to be the current in the large glow discharge of the free space spectrometer, which in the first two surveys was set at 75–100 mA, the current which yields the strongest lines of C_3N , but for C_3N^- turned out to be too high

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TABLE 1
LABORATORY FREQUENCIES OF C_3N^-

Transition	Frequency ^a (MHz)	$O-C$ ^b (kHz)
10–9	97,029.656	–37
12–11	116,434.180	–3
15–14	145,539.368	–27
16–15	155,240.672	11
19–18	184,342.811	0
23–22	223,141.256	34
25–24	242,538.225	3
27–26	261,933.577	2
28–27	271,630.585	–8
29–28	281,327.147	–4
35–34	339,495.912	19
39–38	378,263.734	–16

^a The estimated 1σ measurement uncertainties are 15 kHz.

^b Calculated from the spectroscopic constants in Table 2.

by a factor of 4–5. We were undeterred by our initial failure, because C_3N^- appeared to be an extremely good candidate for both laboratory and astronomical detection: its dipole moment is calculated to be quite large, 3.1 D (this work; Kolos et al. 2008), and its electron affinity (the electron binding energy) is one of the highest of any polyatomic molecule, 4.59 ± 0.25 eV according to Graupner et al. (2006).

We first searched for lines of C_3N^- between 329 and 359 GHz in a 100 mA discharge through a mixture of commercially available gases that yielded an optimum concentration of C_3N : cyanogen (NCCN, 25%), acetylene (C_2H_2 , 25%), and nitrogen (N_2 , 50%) with the walls of the discharge cell cooled to 200 K. Yielding no evidence for C_3N^- , the commercial gases were replaced with cyanoacetylene (HC_3N , synthesized by the procedure described in Appendix A), because it had been shown that HC_3N is the best source of the C_3N radical (Gottlieb et al. 1983; McCarthy et al. 1995) and the C_3N^- anion is produced by dissociative electron attachment to HC_3N at very low electron energies (1.6 eV; Graupner et al. 2006). In the second search, three transitions were covered in a fairly high current discharge (75 mA) through HC_3N (20%) and Ar (80%) cooled to 200 K; again, no lines attributable to C_3N^- appeared. Lines of C_3N^- were finally observed in the third search, which was done with the same mixture of gases as for the previous one, but at a lower current (20 mA). The total pressure was about 15 mtorr when the walls of the discharge cell were cooled to 250 K.

TABLE 2
SPECTROSCOPIC CONSTANTS OF C_3N^-

CONSTANT	LABORATORY	ASTRONOMICAL	THEORETICAL	
			This Work ^a	KGB ^b
B	4851.62183(20)	4851.61(3)	4848	4850
$10^6 D$	685.92(10)	700(100)	627	628
eQq	–3.3	...

NOTE.—Values are given in MHz.

^a B_e (4852.7 MHz) from a CCSD(T)/aug-cc-pVQZ calculation; the vibration-rotation correction (4.6 MHz) and D were calculated at the CCSD(T)/cc-pVTZ level of theory.

^b B_e (4850.6 MHz) and D from a CCSD(T)/aug-cc-pCV5Z calculation; the vibration-rotation correction (0.6 MHz) was calculated at the CCSD(T)/aug-cc-pVQZ level of theory (Kolos et al. 2008).

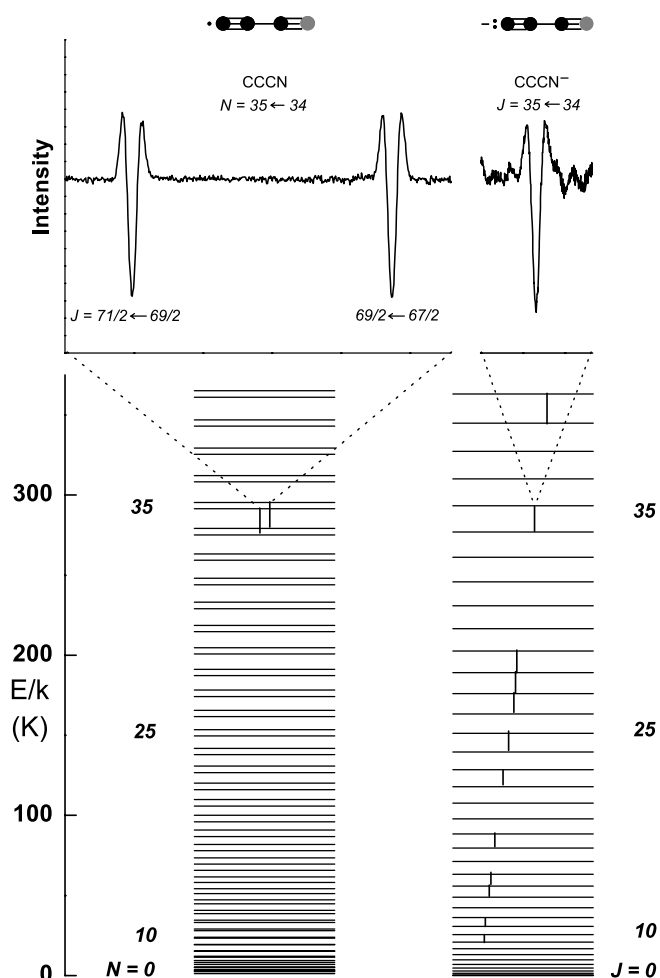


FIG. 1.—Sample laboratory spectra and rotational energy levels of C_3N and C_3N^- , illustrating the collapse of the spin doublet fine structure of C_3N on electron attachment. *Left*: $N = 35 \leftarrow 34$ transition of C_3N observed with an integration time of 3 minutes, showing the well-resolved spin rotation doublets. The separation of the spin doublets are magnified 6000-fold in the energy level diagram. Simultaneous frequency and Zeeman modulation were used to remove background lines; owing to the modulation employed the instrumental line shape is approximately the second derivative of a Lorentzian. *Right*: Corresponding transition of C_3N^- observed in the same discharge with an integration time of 10 minutes. The intensity scale is magnified with respect to that of C_3N . The measured transitions are indicated by solid vertical lines.

The density of background lines in the HC_3N discharge was quite high. Typically one line every 10 MHz was encountered at an intensity comparable to that of C_3N^- , or about 50 background lines at each transition surveyed. The lines of C_3N^- could be distinguished from this clutter of unrelated lines because transitions in the closed-shell $^1\Sigma$ electronic ground state are closely harmonic in frequency, and the theoretical centrifugal distortion constant (D) is sufficiently accurate to predict successive transitions to within the pressure broadened line width (1.2 MHz). Two successive harmonically related lines at 262 and 272 GHz, assigned to the $J = 27-26$ and $28-27$ transitions, yielded a preliminary estimate of the rotational constant B . Additional lines with comparable intensities were then observed at both higher and lower frequencies, confirming the assignment.

The highest concentration of C_3N^- was observed with a discharge current (20 mA) that was 5 times lower than that which was best for C_3N (100 mA). The line at 339 GHz was then about 100 times weaker than the corresponding transition of C_3N (summed over both spin components; see Fig. 1). Owing to the smaller

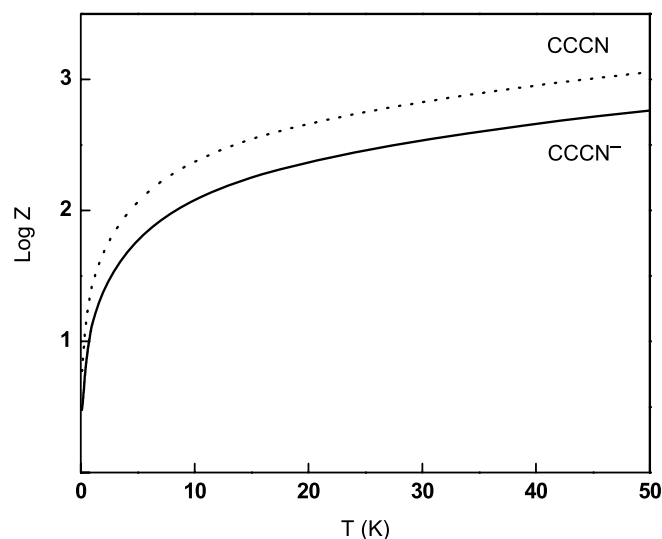


Fig. 2.—Rotational partition function (Z) of C_3N and C_3N^- vs. temperature (T).

partition function and larger dipole moment, lines of C_3N^- are 2.4 times more intense than those of C_3N per unit abundance (Fig. 2). The concentration of C_3N^- in our discharge ($3 \times 10^7 \text{ cm}^{-3}$), corresponding to a mole fraction of 5×10^{-8} , is $\leq 0.001\%$ that of HC_3N under the conditions in which lines of the anion are most intense. The anion to radical concentration—about 2%—is comparable to that of CN^- (Gottlieb et al. 2007), but the mole fraction of C_3N^- is about 100 times smaller than that of CN^- in the respective discharges. Following detection of C_3N^- in a discharge through HC_3N (with either Ar or N_2 as the carrier gas), other known sources of C_3N were examined for evidence of the anion. Lines of C_3N^- were observed in low-current (20 mA) discharges through $HCCH/N_2$ and $HCCH/N_2/NCCN$, but not through $NCCN$ with either Ar or N_2 as the carrier. Although the latter discharge is a good source of the radical, the anion was not observed because of the absence of hydrogen.

The evidence that C_3N^- is the carrier of the series of lines in Table 1 is overwhelming. As mentioned, the assigned lines are closely harmonic in frequency, and show no evidence of fine or hyperfine structure, as expected for high rotational transitions of a molecule in a $^1\Sigma$ ground electronic state. The rotational constant (B) is only 2% smaller than that of C_3N , about as expected, because electron attachment to carbon chains generally results in only a slight decrease in B . There is no evidence for nitrogen quadrupole hyperfine structure and none is expected, because hyperfine structure (hfs; $< 10 \text{ kHz}$ at 97 GHz) is much less than the pressure broadened line width. Although we did not directly rule out H as a possible constituent by observing the same series of lines in a discharge through DC_3N , the addition of an H atom at the end of the chain would yield a rotational constant that is too small; for example, B for HC_3N is 6% smaller than that of the carrier of the new series of lines.

The spectroscopic constants provide strong additional evidence in support of the assignment. First, the measured rotational constant is within 0.1% of that predicted by two independent high-level quantum calculations, and the centrifugal distortion (D) is in excellent agreement with that calculated. Second, other possible carriers such as the carbon-13 or nitrogen-15 isotopic species of the C_3N radical are ruled out because: (1) the rotational spectra of these have been measured (McCarthy et al. 1995); and (2) the spectra of the isotopic species of the radical would consist of symmetrical spin doublets separated by about 19 MHz, rather than the single series of lines observed. Third, vibrationally excited

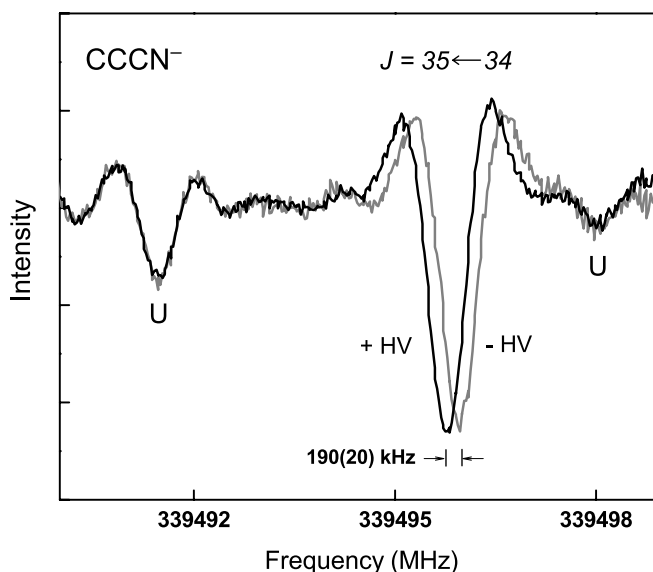


Fig. 3.—Ion drift measurement of C_3N^- . The $J = 35 \leftarrow 34$ absorption line of C_3N^- with either positive or negative high voltage (HV) applied to the electrode near the radiation source; the electrode near the detector was held at ground. The integration time for each spectrum was approximately 30 minutes. The frequency shift derived here ($95 \pm 10 \text{ kHz}$) corresponds to an axial ion drift velocity of $80 \pm 10 \text{ m s}^{-1}$. The weak features marked U are from unassigned neutral molecules produced in the discharge.

C_3N is ruled out for several reasons: the rotational constants in the low-lying bending vibrational levels are close to that of the ground state (Mikami et al. 1989); although the rotational spectra in the stretching levels of C_3N have not been measured, by analogy with other linear polyatomic molecules such as HC_3N , the rotational constants are expected to be within 1% of that of the ground state; and the symmetry of the lowest stretching levels in C_3N is $^2\Sigma$, so the spectra will consist of well-resolved spin doublets rather than a single series of lines. Fourth, the C_3N^+ cation is ruled out because three published quantum calculations predict it to have a triplet electronic ground state (Harland & MacLagan 1987; Ding et al. 1998; Kolos et al. 2008). The rotational spectrum of C_3N^+ would then consist of three well-resolved fine-structure components of comparable intensity, rather than the observed single series of harmonic lines.

Finally, direct determination of the charge of the carrier provides conclusive confirmation of our assignment. Ions are accelerated in the positive column of a glow discharge owing to the axial electric field [typically $(5\text{--}15) \text{ V cm}^{-1}$; Stephenson & Saykally 2005]. A measurable Doppler shift in the line frequency was observed when the millimeter-wave radiation made a single pass through the discharge cell (normally a double pass is employed to enhance sensitivity). The sign of the shift then changed on reversal of the cathode and anode potentials, as expected for a negative ion (Fig. 3). The observed ion drift velocity ($v_{da} = 80 \pm 10 \text{ m s}^{-1}$) is within 10% of that inferred from earlier measurements of CN^- in a similar discharge (Gottlieb et al. 2007) on the assumption that v_{da} is inversely proportional to the square root of the mass of the ion.

2.2. Astronomical

2.2.1. IRC+10216

The recent detection of C_4H^- in IRC+10216 with the IRAM 30 m telescope implies that the isoelectronic C_3N^- might be detectable in this source. Before undertaking new observations, we examined spectra of this source obtained in a deep survey for

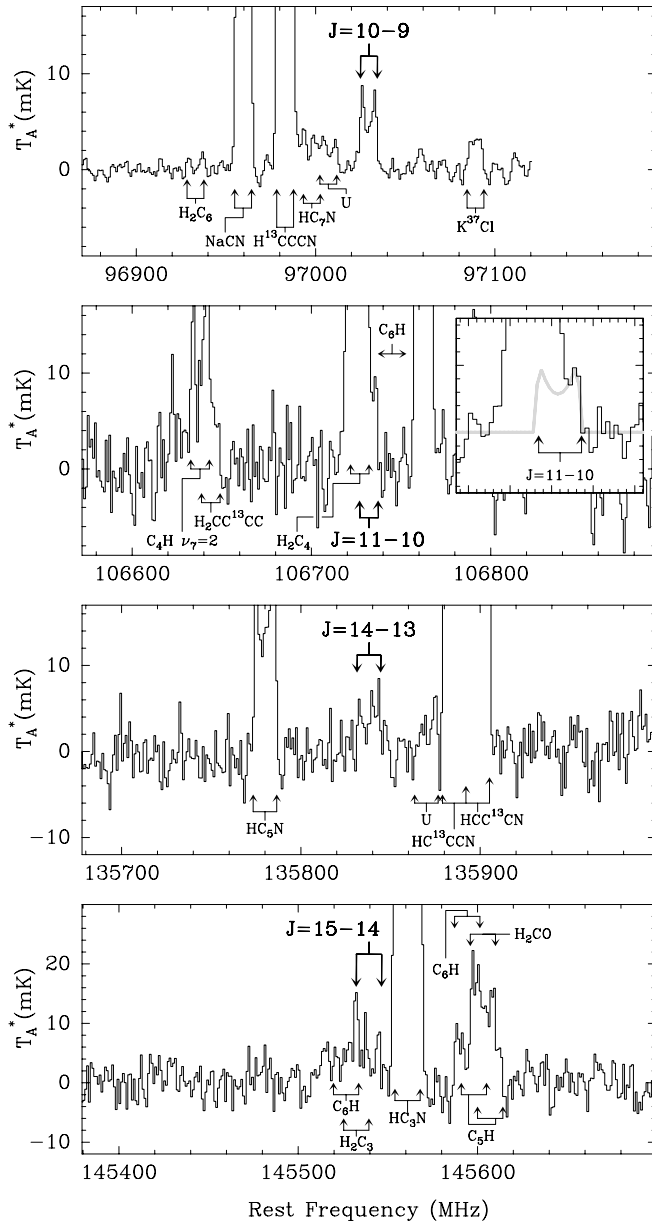


FIG. 4.—Four transitions of C_3N^- observed toward IRC+10216. The red and blue edges of the cusped line shapes are indicated by vertical arrows. The $J = 11-10$ line at 106.7 GHz is blended with an intense line of C_6H (see simulated profile in insert), and the $J = 15-14$ line at 145.5 GHz is partially blended with very weak lines of $l-H_2C_3$ and C_6H . The antenna temperature scale is corrected for atmosphere attenuation and spillover losses (T_A^*), the spectral resolution is 1 MHz, and the frequency axis is relative to a systemic velocity of -26.5 km s^{-1} .

new molecules in the 3 mm band (J. Cernicharo et al. 2008, in preparation). This survey contains a number of lines that could not be assigned to any known carrier, including two of similar intensity which agree in frequency to better than 0.5 MHz with the $J = 10-9$ (97,030 MHz) and $11-10$ (106,732 MHz) transitions of C_3N^- . The evidence for C_3N^- is compelling, because the density of lines in our 3 mm spectral survey with intensities $\geq 5 \text{ mK}$ is only about one every 30 MHz, and the number of lines with $T_A^* \sim 7 \text{ mK}$ is even less: roughly one every 100 MHz. Although three other transitions of C_3N^- were also covered in this band (at 77,625, 87,327, and 116,434 MHz), these are either blended with intense background lines [e.g., the $(N, J, F) = (1, 3/2, 1-0, 1/2, 0)$ component of C_2H at 87,328 MHz], or they lie at frequencies where the noise temperatures of the receivers

TABLE 3
OBSERVED LINE PARAMETERS OF C_3N^- IN IRC+10216

TRANSITION $J'-J$	FREQUENCY (MHz)		$\int T_A^* dv$ (K km s $^{-1}$)	V_{exp} (km s $^{-1}$)
	Calculated ^a	Observed		
10–9	97029.693	97029.5(3)	0.183(12)	14.1(3)
11–10	106732.028	106731.5(8)	0.224(20) ^b	14.5 ^c
14–13	135837.883	135837.9(6)	0.128(20)	14.1(6)
15–14	145539.395	145538.6(10)	0.165(40) ^b	14.5 ^c

NOTES.—The numbers in parentheses are the 1σ uncertainties in units of the last digits. The observed line frequencies are computed for a source velocity $v_{\text{LSR}} = -26.5 \text{ km s}^{-1}$.

^a Calculated with the spectroscopic constants derived from the laboratory measurements (Table 2).

^b Partially blended line.

^c Terminal velocity constrained to 14.5 km s^{-1} .

are high. To confirm our tentative identification of C_3N^- , we then undertook observations in 2007 September and October of two additional transitions in the 2 mm wave band.

The observations in the 2 mm band were also made with the 30 m telescope. The dual-polarization receivers were operated in the single-sideband mode with a rejection of the upper sideband of about 18 dB, as estimated from the intensities of strong lines. The back ends consisted of two 512 channel filter banks with 1 MHz resolution. To ensure flat baselines, the secondary mirror was wobbled by $\pm 90''$ at a 0.5 Hz rate. The zenith sky opacity of 0.03–0.20 resulted in a system noise temperature of 170–330 K. The total integration time for each polarization in the position switching mode was 4 hr at 135.8 GHz and 2 hr at 145.5 GHz, which yielded an rms noise of 2–3 mK when the two polarizations were averaged.

Two lines of similar intensity ($T_A^* \sim 6 \text{ mK}$) coincide in frequency with the $J = 14-13$ (135,838 MHz) and $15-14$ (145,539 MHz) transitions of C_3N^- . Although the line at 145.5 GHz is partially blended with lines of C_6H and $l-H_2C_3$, it could be assigned to C_3N^- with some confidence, because the line width (29 km s $^{-1}$) and cusped line shape—characteristic of molecular emission from the outer envelope of IRC+10216—are similar to that of the line at 97 GHz, which was observed with excellent signal to noise and is free of contamination by background lines. All four lines are shown in Figure 4; the data are summarized in Table 3.

Figure 5 shows the rotational temperature diagrams for lines of C_3N^- and C_3N observed in the IRAM 30 m spectral survey (Cernicharo et al. 2008, in preparation, 2000; Table 3, this work). The radius of the shell as determined from that of the neutral species is $10''$ – $24''$ (see, e.g., Guélin et al. 1993), so line intensities have been corrected for beam dilution assuming a uniform source of radius $20''$. The rotational temperature T_{rot} is the same for the anion and the radical to within the measurement uncertainties: $24 \pm 5 \text{ K}$ for C_3N^- and $29 \pm 2 \text{ K}$ for C_3N . The total column density across the shell is $(1.6 \pm 0.6) \times 10^{12} \text{ cm}^{-2}$ for C_3N^- versus $(3.1 \pm 0.3) \times 10^{14} \text{ cm}^{-2}$ for C_3N . The abundance ratio $[C_3N^-]/[C_3N]$ of 0.52% is about 20 times larger than $[C_4H^-]/[C_4H]$, but is 12 times smaller than $[C_6H^-]/[C_6H]$ (Cernicharo et al. 2007) and 50–70 times smaller than $[C_8H^-]/[C_8H]$ in this source (Remijan et al. 2007; Kawaguchi et al. 2007).

There is little doubt that C_3N^- is the carrier of the four assigned lines in IRC+10216. As expected, the intensities of all four are similar, and there are no missing lines. The spectroscopic constants B and D are in excellent agreement with those from the laboratory (Table 2). The rotational temperature is the same as that of C_3N , as expected (Fig. 5). The line observed with

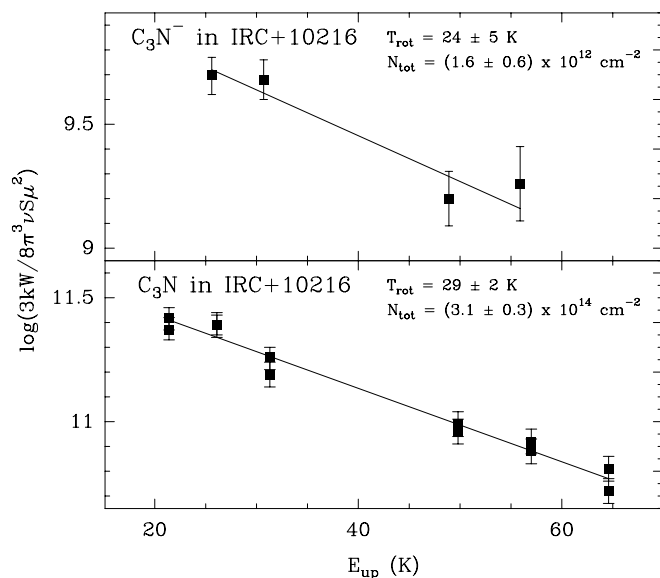


FIG. 5.—Rotational temperature diagram for C_3N^- and C_3N in IRC+10216. The measured line intensities were corrected for beam dilution assuming a uniformly bright source of radius $20''$.

the highest signal-to-noise ratio ($J = 10-9$) is strongly cusped, confirming that C_3N^- is confined to the same hollow shell as C_3N and the neutral acetylenic chain radicals, e.g., C_3H , C_4H , and C_6H .

2.2.2. TMC-1

The recent observations of C_6H^- and C_8H^- in TMC-1 with abundances of a few percent that of the corresponding radical (McCarthy et al. 2006; Brünken et al. 2007), suggested that C_3N^- might be detected in this source. In 2007 October we did a simul-

taneous search with the NRAO 100 m Green Bank Telescope (GBT) for the 2–1 lines of C_3N^- and C_4H^- in TMC-1 near 19 GHz, and measurements of the corresponding lines of C_3N and C_4H . System temperatures with the K-band receiver were typically 30–50 K. Spectra were acquired as before by position switching (Brünken et al. 2007), at a resolution of 6.1 kHz across the 50 MHz wide bands of the spectrometer.

Column densities were derived from the line parameters in Table 4 on the assumption that the rotational temperature is 5 K. The column density of C_3N^- in TMC-1 of $\leq 7 \times 10^{11} \text{ cm}^{-2}$ is $\leq 0.8\%$ of C_3N ($9.5 \times 10^{13} \text{ cm}^{-2}$). Our limit for the $\text{C}_3\text{N}^-/\text{C}_3\text{N}$ ratio in TMC-1 is comparable to the ratio in IRC+10216 (0.52%; § 2.2.1). For C_4H^- our limit of the column density in TMC-1 ($\leq 2.3 \times 10^{10} \text{ cm}^{-2}$) is nearly 4 times lower than a preliminary estimate obtained by Brünken et al., and is 2 times lower than that determined in a search for the 83.8 GHz line (Agúndez et al. 2008). The $\text{C}_4\text{H}^-/\text{C}_4\text{H}$ ratio in TMC-1 ($< 0.004\%$ for a C_4H column density⁵ of $5.7 \times 10^{14} \text{ cm}^{-2}$) is 5 times lower than that found in IRC+10216 (0.02%; Cernicharo et al. 2007). The upper limit for the anion to radical ratio for C_4H is much lower than for C_3N , because lines of C_4H are several times more intense than those of C_3N in TMC-1 and the spectroscopic enhancement for C_4H^- (100; Gupta et al. 2007) is about 40 times greater than for C_3N^- (2.4; see Appendix B).

3. DISCUSSION

Our laboratory search for C_3N^- was guided by high-level quantum calculations of the structure of this anion. The equilibrium geometry was obtained at the coupled cluster [CCSD(T)] level of theory with a large basis set. The anharmonic force field calculated by the method described by Stanton & Gauss (2000) was

⁵ There is a typographical error for the C_4H column density in Table 2 of Brünken et al. (2007). The correct value is $6.1 \times 10^{14} \text{ cm}^{-2}$.

TABLE 4
OBSERVED LINE PARAMETERS OF C_3N AND C_4H , AND UPPER LIMITS OF C_3N^- AND C_4H^- IN TMC-1

MOLECULE	TRANSITION			FREQUENCY ^a (MHz)	T_A^* (mK)	V_{LSR} (km s ⁻¹)	Δv (km s ⁻¹)
	$N'-N$	$J'-J$	$F'-F$				
C_3N^-	1–0	...	2–1	9,703.406	<21
	2–1	...	3–2	19,406.536	<12
C_3N	2–1	5/2–3/2	5/2–3/2	19,780.808	203(9)	5.72	0.25(1)
	3/2–1/2	19,780.823	128(9)	5.54	0.21(2)
	5/2–5/2	19,781.005	47(7)	5.83	0.35(6)
	7/2–5/2	19,781.090	344(9)	5.56	0.24(7)
	3/2–3/2	19,783.537	56(10)	5.62	0.16(3)
	...	3/2–1/2	3/2–3/2	19,797.472	59(8)	5.62	0.23(4)
	3/2–1/2	19,799.959	71(7)	5.69	0.26(3)
	5/2–3/2	19,800.126	205(8)	5.69	0.25(1)
	1/2–1/2	19,800.971	65(7)	5.65	0.27(3)
C_4H^-	1–0	9,309.887	<8
	2–1	18,619.758	<7
C_4H	1–0	3/2–1/2	1–0	9,493.060	186(7)	5.67	0.33(2)
	2–1	9,497.615	605(11)	5.64	0.31(7)
	1–1	9,508.005	162(7)	5.61	0.33(2)
	2–1	5/2–3/2	2–1	19,014.720	1033(44)	5.72	0.26(1)
	3–2	19,015.144	1527(45)	5.74	0.26(1)
	2–2	19,025.107	105(20)	5.72	0.46(10)
	...	3/2–1/2	1–1	19,044.760	162(22)	5.75	0.36(6)
	2–1	19,054.476	1010(30)	5.74	0.27(1)
	3–2	19,055.947	440(30)	5.72	0.23(2)

NOTE.—Position assumed for TMC-1: $\alpha(\text{J2000.0}) = 04^{\text{h}}41^{\text{m}}42.49^{\text{s}}$, $\delta(\text{J2000.0}) = +25^{\circ}41'26.9''$.

^a Frequencies of C_3N from unpublished measurements in a supersonic molecular beam (estimated 1 σ uncertainties 5 kHz), C_4H frequencies from Gottlieb et al. (1983), and C_4H^- from Gupta et al. (2007).

then used to calculate the vibration-rotation interaction constants (α_i), and to convert the equilibrium rotational constant B_e to that of the ground state $B_0 = B_e - \frac{1}{2} \sum_i \alpha_i$. After our calculation was completed, we learned of an independent calculation of B_0 at the same level of theory but with a larger basis set (Kolos et al. 2008). As Table 2 shows, the two calculations agree to better than 0.1% with the measured B_0 .

The formation of anions in our glow discharge is poorly understood. In our earlier work on CN^- , we discussed evidence for two mechanisms: (1) radiative electron attachment to the CN radical, and (2) dissociative attachment to the precursor $NCCN$. A significant limitation in the interpretation of the CN^- observations is that the ratio of CN^- to $NCCN$ cannot be determined, owing to the nonpolarity of $NCCN$. In the experiments here, the relative concentrations of C_3N^- and HC_3N were determined from the measured line intensities under conditions that yielded the most intense lines of the anion. From these measurements, it appears that C_3N^- may be formed by dissociative attachment to HC_3N , because the C_3N^-/HC_3N ratio $[(3-10) \times 10^{-6}]$ is similar in discharges through different precursor gases, and C_3N^- was not observed in a discharge through $NCCN$ and N_2 in which C_3N was produced but not HC_3N .

We have attempted without success to detect the lowest rotational transitions of C_3N^- with our Fourier transform microwave (FTM) spectrometer employing a variety of experimental conditions, and both cyanoacetylene and cyanogen as precursor gases. With either gas mixture, lines of C_3N radical are sufficiently strong that transitions of its carbon-13 isotopic species (McCarthy et al. 2003) are readily observed in natural abundance, but there is no evidence for C_3N^- at a level of 10^{-3} of the neutral radical.

This failure to find C_3N^- is puzzling; it may mean that anion production in the supersonic beam of the FTM is via radiative electron attachment, and that the efficiency of electron capture is significantly enhanced for radicals with large dipole moments. Lines of isoelectronic C_4H^- , for example, were only observed with He or H_2 as a buffer gas (Gupta et al. 2007), even though lines of ground state C_4H were more than 10 times weaker in He or H_2 than in Ne or Ar. Detection of C_4H^- under these conditions may occur because electron attachment is not from the ground state, which possesses a small dipole moment (0.87 D), but rather from the low-lying $^2\Pi$ state (0.01 eV) which has a large dipole moment (4.5 D; see Pachkov et al. 2003) and is thermally populated in the warmer He or H_2 expansions, but is not populated in the Ne or Ar expansions. Although C_3N has an even higher electron affinity (4.59 eV) than C_4H , the dipole moment of the ground state (2.9 D) is only about half that of the $^2\Pi$ electronic state of C_4H , and its first excited electronic state (0.2 eV) is only weakly polar (0.2 D; McCarthy et al. 1995). The apparent absence of C_3N^- lines in the FTM beam may mean that the dipole moment plays a crucial role in determining the rate of electron attachment.

As predicted (e.g., Millar et al. 2000) the abundances of C_6H^- and C_8H^- in TMC-1 and IRC+10216, and C_6H^- in L1527 are fairly large, representing a good fraction of that of their parent neutral species ($\sim 10\%$). On the other hand, the abundance of C_4H^- relative to C_4H is much smaller (0.025% in IRC+10216, 0.01% in L1527, and $<0.004\%$ in TMC-1; Cernicharo et al. 2007; Agúndez et al. 2008; see § 2.2.2). This may simply be because C_4H is smaller than C_6H and C_8H and has far fewer low-lying vibrational states, making electron attachment less probable. The electron attachment rate coefficients, like all two-body radiative attachment rate coefficients, are difficult to calculate precisely, and are generally evaluated via a statistical theory (see Herbst

1980). Such a theory seems to work for large carbon chains, where nearly every collision between a neutral and an electron yields an anion, but seems to fail for shorter chains: the rate coefficient for $C_4H + e$ calculated by E. Herbst (2007, private communication) deviates significantly from that derived from the observed C_4H^- abundance (see Cernicharo et al. 2007). Further testing of the statistical theory on other short anions is crucial, particularly because this theory is also used to calculate the rates of radiative association between atoms and neutrals, important in astrochemistry.

A crucial parameter that may control the abundance of molecular anions in the interstellar gas is the overall rate of radiative attachment (k_{ratt}) for the process $A + e \rightarrow A^- + h\nu$. Following Cernicharo et al. (2007), we derived $k_{\text{ratt}}(C_3N)$ from the ratio of the observed abundances of C_3N and C_3N^- in the expanding shell of IRC+10216. On the assumption that the steady state approximation applies and photodetachment may be neglected,

$$k_{\text{ratt}} = \frac{[A^-]}{[A]} \frac{1}{[e]} (k_{\text{aa}}[H] + k_{\text{neu}}[B^+]), \quad (1)$$

where $[e]$ is the abundance of electrons, $[B^+]$ of cations, and $[H]$ of atomic hydrogen; k_{aa} is the rate coefficient for associative detachment by reaction with H; and k_{neu} for neutralization with cations. We adopt $k_{\text{neu}} = 10^{-7} \text{ cm}^3 \text{ s}^{-1}$ (see, e.g., Smith & Adams 1979) and k_{aa} from the measurements of Barckholtz et al. (2001). The abundances of the cations ($[e]/[H_2] = [B^+]/[H_2] = 10^{-7}$; Glassgold et al. 1987) and atomic hydrogen ($[H]/[H_2] = 10^{-4}$; Glassgold 1996) were estimated for a shell of radius $15''$.

From the observed values $[C_4H^-]/[C_4H] = 1/4200$ and $[C_6H^-]/[C_6H] = 1/16$ (Cernicharo et al. 2007), we arrive at $k_{\text{ratt}}(C_4H) = 2 \times 10^{-10} \text{ cm}^3 \text{ s}^{-1}$ and $k_{\text{ratt}}(C_6H) = 5 \times 10^{-8} \text{ cm}^3 \text{ s}^{-1}$. A more thorough time-dependent calculation involving expansion of the envelope and the loss of anions through reactions with neutral atoms other than H and by photodetachment yields $k_{\text{ratt}}(C_4H) = 9 \times 10^{-11} (T/300)^{-0.5} \text{ cm}^3 \text{ s}^{-1}$, $k_{\text{ratt}}(C_6H) = 1.4 \times 10^{-8} (T/300)^{-0.5} \text{ cm}^3 \text{ s}^{-1}$, and $k_{\text{ratt}}(C_8H) = 2.5 \times 10^{-8} (T/300)^{-0.5} \text{ cm}^3 \text{ s}^{-1}$ (see Agúndez et al. 2008). Similarly, from the observed values of $[C_3N^-]/[C_3N]$, we derive $k_{\text{ratt}}(C_3N) = 1.9 \times 10^{-9} (T/300)^{-0.5} \text{ cm}^3 \text{ s}^{-1}$ from our time-dependent model calculations. The electron attachment rate coefficient derived for C_3N is then ~ 20 times larger than that of C_4H .

This value of $k_{\text{ratt}}(C_3N)$ may be compared with that derived by Petrie & Herbst (1997), who calculated the total energies and enthalpies of formation of C_3N^- , C_3N , and $CCCNH$; and the intensities and vibrational frequencies of C_3N^- . They derived $k_{\text{ratt}}(C_3N) = (2-4) \times 10^{-10} (T/300)^{-0.5} \text{ cm}^3 \text{ s}^{-1}$, depending on whether they used the calculated or the experimental electron affinity of C_3N . This is almost an order of magnitude smaller than the apparent rate derived from our measured C_3N^- abundance.

Another possible pathway for the formation of C_3N^- is dissociative attachment to HNCCC (a metastable isomer of cyanoacetylene) first proposed by Petrie & Herbst (1997),



The column density of HNCCC derived from recent observations of the $J = 9-8$ through $12-11$ lines in IRC+10216 with the IRAM telescope, $7 \times 10^{11} \text{ cm}^{-2}$, is about 500 times less than that of C_3N (J. Cernicharo et al. 2008, in preparation). Therefore, reaction (2) would have to be at least 500 times faster than reaction (1) to account for the observed $[C_3N^-]/[C_3N]$ ratio in IRC+10216.

Such a high rate of $\geq 10^{-6} \text{ cm}^3 \text{ s}^{-1}$ for reaction (2) seems implausible. HNCCC is relatively abundant in TMC-1 (Snyder et al. 2006) but C_3N^- is not detected there, implying that reaction (2) may not be a significant channel for the formation of C_3N^- in cold clouds. If reaction (1) is the only efficient path for the formation of C_3N^- in the interstellar gas, then $k_{\text{ratt}}(\text{C}_3\text{N})$ derived from the measured abundance of C_3N^- in IRC+10216 is the true rate of radiative attachment, which we have shown is about an order of magnitude greater than that predicted by Petrie & Herbst (1997).

The laboratory identification of new molecular ions in the millimeter-wave band is much easier when it is possible to readily distinguish ionic lines from those of neutral molecules. The technique of velocity modulation provides excellent discrimination between ions and neutrals in the IR band (Stephenson & Saykally 2005), but when lines of the ions are observed by this technique in the millimeter-wave band they are about 6 times weaker than those observed in a normal DC glow discharge (because the Doppler shift from the ion drift is much less than the pressure broadened line width; Savage & Ziurys 2005; see also Fig. 3). If the sensitivity of our millimeter-wave spectrometer can be increased moderately, say by a factor of 5, velocity modulation may become a powerful general technique for the detection of new molecular anions.

Other anions of astrophysical interest may be detected by the present technique. Five good candidates are CH_3CC^- (1-propynyl), $\text{c-C}_3\text{H}^-$, C_5N^- , H_2CCC^- , and H_2CCN^- . These span a wide range of binding energy and dipole moment: linear C_5N^- and the symmetric top CH_3CC^- have large dipole moments and binding energies that are comparable to that of C_6H^- , but the binding energies of the others are ≤ 2.0 eV. The rotational constants and leading centrifugal distortion constants of H_2CCC^- and H_2CCN^- were derived by autodetachment spectroscopy (Yokoyama et al. 1996; Lykke et al. 1987), so the search for the rotational lines can be restricted to a fairly narrow frequency range. Observations of these five molecular anions may allow astronomers to better

determine the electron density in the interstellar gas, when the dependence of the cross sections for electron attachment and photodetachment on dipole moment and electron affinity have been determined.

For new observations, the highest priority in our view should be given to the discovery of better astronomical sources for negative molecular ions in a wide range of objects, starting with a deep survey of C_6H^- , the most readily observed molecular anion to date. It is highly implausible that with IRC+10216, TMC-1, and recently L1527 the best anion sources in the entire Galaxy have been found; it is more reasonable to assume instead that this limitation is largely the well-known selection effect of observers devoting their limited observing time on big telescopes to the few sources which have worked in the past. The same selection effect was encountered in the early study of interstellar molecules. Breaking free of this straitjacket is worth a substantial amount of observing time on the largest radio telescopes. It is not premature to suggest that such a survey should be extended to some of the more plausible extragalactic sources.

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APPENDIX A

Cyanoacetylene (HC_3N) was synthesized 20 g at a time with a yield of about 70% by a procedure similar to that described by Miller & Lemon (1967). To ensure a high yield of propiolamide (HCCCONH_2 , the product of the first step in the synthesis), the large excess of NH_3 was removed by flushing the reaction vessel with N_2 and bubbling the NH_3/N_2 mixture through concentrated HCl. Methanol (CH_3OH) present as a by-product in the reaction was then pumped away. The sample was stored indefinitely at -20°C with no discernible evidence of polymerization, but during the experiment HC_3N was maintained at a slightly higher temperature (-10°C) to provide sufficient vapor pressure (~ 60 torr; Dannhauser & Flueckinger 1963). A flow rate of $3\text{--}4 \text{ cm}^3 \text{ minute}^{-1}$ (estimated with a mass flow controller), yielded near optimum line intensities of the C_3N radical and a reasonable rate of HC_3N consumption of approximately 0.5 g hr^{-1} . In most experiments described here the flow of HC_3N vapor was maintained with a manual metering valve because it was found that in a few days running the electronic flow controller becomes clogged. As a result, the flow rate of HC_3N is less certain than that of the other precursor gases, owing to changes in the vapor pressure due to variations in the temperature of the sample bath.

APPENDIX B

Shown in Figure 1 are sample spectra and rotational energy level diagrams of C_3N and C_3N^- . The spin rotation doublet splittings in the energy level diagram of C_3N are magnified to illustrate the higher density of levels in the $^2\Sigma$ electronic ground state of the radical. Electron attachment to the radical results in the collapse of the spin doublets, simplifying the spectrum of the anion to that of closed shell linear molecule in a $^1\Sigma$ ground state. The enhancement of the line intensities of C_3N^- relative to those of C_3N by a factor of 2.4 results mainly from the collapse of the spin doublets of the radical, since the dipole moment of the anion (3.1 D) is only slightly larger than that of the radical (2.85 D; McCarthy et al. 1995). As Figure 2 shows, this spectroscopic enhancement of C_3N^- relative to C_3N is independent of temperature for the typical astronomical sources where molecular anions have been observed.

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Note added in proof.—After submission of this paper, we observed the five lowest rotational transitions of C_3N^- by means of Fourier transform microwave (FTM) spectroscopy. These new measurements provide frequencies of the hyperfine-split components which are accurate to better than 0.2 ppm, i.e., a fraction of a linewidth in TMC-1. The nitrogen quadrupole coupling constant derived here, $eQq = -3.248 \pm 0.005$ MHz, is in excellent agreement with the theoretical prediction of -3.3 MHz. A complete account of the FTM measurements will be presented elsewhere.